



# Implicit Solvation Models in ONETEP

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#### Importance of solvation

- Many important biochemical reactions occur in aqueous solution.
- Performing calculations *in vacuo* often leads to greatly inaccurate results. Especially sensitive properties and phenomena include, among others:
  - energy differences between molecular conformers [1],
  - rates of reactions [2],
  - tautomeric equilibria [1],
  - $-\pi$ -facial sensitivity [2],
  - molecular (esp. protein-protein) associations [1,3],
  - protein structures [4],
  - ligand binding free energies [5].
- Thus proper description of the solvent environment is crucial in simulations of biological molecules.

### Explicit solvent approach

- Introduce the solvent in molecular detail.
- Pros:
  - provides accurate treatment of solute-solvent interactions,
- Cons:
  - leads to an increase in system size, possibly by an order of magnitude [6],

Animation by C. Pittock



- must average out instantaneous interactions before the results become meaningful, ("integrate out" the degrees of freedom of the solvent) [4],
- to avoid surface effects, PBCs must be introduced. This in turn requires large cells to avoid artificial interactions between replicated solutes.

#### Implicit solvent approach

- Treat the solute explicitly and embed it within a suitably defined cavity, the inside of which is inaccessible to the solvent.
- Replace the solvent with an *unstructured dielectric continuum,* only retaining its **average** effect on the solute [7].
- The charge distribution of the solute polarizes the dielectric, creating a reflection charge, which then interacts with the solute [2].



#### Implicit solvent approach

- In other words: one assumes that the macroscopic description of the solvent as continuous dielectric medium can be used as an approximation on the microscopic scale [8].
- This works because the largest part of molecular interactions in solution is characterized by low specificity and low directionality [9].
- Pros:
  - no solvent atoms necessary,
  - eliminates the costly sampling of the solvent degrees of freedom [6],
  - faster sampling of solute conformations owing to absence of solvent friction [5].



Effect of solvent on the potential, for CN<sup>-</sup>



Potential, U [V]

Effect of solvent on the potential, for CN<sup>-</sup>



#### Two terms in free energy of solvation

only the total can

be obtained 
$$\longrightarrow \Delta G_{sol} = \Delta G_{pol} + \Delta G_{npol}$$

- The **electrostatic** or **polar term** describes the response of the solvent to the charge distribution of the solute [7].
- It is the difference between the electrostatic energy

$$\frac{1}{2}\int \rho(\mathbf{r})\phi(\mathbf{r})d\mathbf{r}$$

in solvent and in vacuum.

# Obtaining $\phi(\mathbf{r})$

• In vacuum the Poisson equation is solved:

$$\varepsilon \nabla^2 \phi(\mathbf{r}) = -4\pi \rho(\mathbf{r})$$

This is easiest to compute in reciprocal space and that's the way it's done in ONETEP.

• With an inhomogeneous dielectric the *generalized* Poisson equation needs to be solved:

 $\nabla \cdot \varepsilon(\mathbf{r}) \nabla \phi(\mathbf{r}) = -4\pi \rho(\mathbf{r})$ This must be solved in real space with a suitable solver.





#### Two terms in solvation energy

 $\Delta G_{\rm sol} = \Delta G_{\rm pol} + \Delta G_{\rm npol}$ 

- The **nonpolar term** accounts for the entropic cost of forming a cavity within the solvent and for the van der Waals interaction of the solute with the solvent [4].
- Difficult to describe rigorously, heuristic approaches are used.
- Most widely used approach is to represent it as a linear function of the molecular surface area [7]:

 $\Delta G_{npol} = \gamma A_{SA}$ .

### Important choices

- How the charges are represented:
  - as classical point charges,
    - VS.
  - distributed (charge densities).
- How solvation is treated:



- for a frozen system: the energy of solvation is determined once, by adding the implicit solvent to the system and observing the resulting change in energy,
- self-consistently: by re-calculating the solvent's influence during every step of energy minimization. This allows the electronic charge density to "feel" the polarizing effect of the solvent and to adapt its shape accordingly.
- Cavity shape.

## Spherical cavity

#### • Pros:

- simple (to implement),
- simple (runs fast).

#### • Cons:

- simple (not especially realistic).
- molecules that are not compact cannot be easily fit into a spherical cavity.
- Has enjoyed considerable success nevertheless ([2] & refs therein).



#### Interlocking spheres

 $\epsilon = 80$ 

 $\epsilon = 1$ 

- Employed in the PCM approach developed by Tomasi, Scrocco and Miertus [9].
- Pros:
  - More realistic cavity.
  - Applicable to molecules of any shape.
- Cons:
  - Requires extra input: parametrization for the sphere radii.

#### Arbitrariness of the cavity



#### **Density-dependent cavity**



#### Density-dependent cavity of Fattebert and Gygi



#### **Density-dependent cavity**



## Implementing IS in ONETEP

- Interface with APBS [10] (Adaptive Poisson-Boltzmann Solver) code.
- APBS is a multigrid solver that can calculate φ(**r**) given ρ(**r**) and ε(**r**). It also automatically calculates the polar term of the solvation energy as:

$$\Delta G_{\text{pol}} = \frac{1}{2} \int \rho(\mathbf{r}) \phi_{\text{E=80}}(\mathbf{r}) d\mathbf{r} - \frac{1}{2} \int \rho(\mathbf{r}) \phi_{\text{E=1}}(\mathbf{r}) d\mathbf{r}$$

• We are currently pursuing both the frozen density approach and the self-consistent one.

## Grid and padding



APBS's box

- Charge densities, potentials are represented on a grid commensurate with ONETEP's *fine grid*.
- Currently using d = 0.136 Å (KE cutoff  $\approx 780 \text{ eV}$ ).
- A small amount of padding ( $\approx 1 \text{ Å}$ ) is necessary.

#### Frozen density approach

#### Implementing IS in ONETEP



Frozen density approach

#### First results



#### First results



#### Correcting the effect of PBCs



Periodic images of the system interact with each other.



[11] Scherlis, Fattebert, Gygi, Cococcioni and Marzari, J. Chem. Phys. **124** (2006).

$$E_{\text{gas}} = E_{\text{PBC}} + c_1 \frac{q^2}{L} + c_2 \frac{qQ}{L^3} + O(L^{-5}) \begin{bmatrix} 2 \\ p \end{bmatrix}$$

[12] Makov and Payne, *Phys. Rev.* B **51** (1995).

#### Frozen density approach

#### First results



#### First results



Frozen density approach

#### Preparation



 An APBS calculation is performed in vacuum for the cores only, to obtain the Coulombic potential due to cores, V<sub>+</sub>(r).



part 1







$$z(\mathbf{r}) = \sum_{A} Z_{A} \widetilde{\delta} \left( \mathbf{r} - \mathbf{R}_{A} \right)$$

charge density of spilled cores

 $n(\mathbf{r})$  charge density of electrons

 $q^{
m die}({f r})$ 

induced charge density (not directly accessible from APBS)



 $V_{+}(\mathbf{r})$ 

Coulombic potential due to the cores

 $V_{-}(\mathbf{r})$ 

Coulombic potential due to the electrons  $V^{
m die}({f r})$ 

potential due induced charge (dielectric's response)

$$z(\mathbf{r}) = \sum_{A} Z_{A} \widetilde{\delta} \left( \mathbf{r} - \mathbf{R}_{A} \right)$$

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1. 
$$E_{\rm e-e} = \frac{1}{2} \int n(\mathbf{r}) V_{-}(\mathbf{r}) d\mathbf{r}$$

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$$E_{\rm e-e} = \frac{1}{2} \int n({\bf r}) V_{-}({\bf r}) d{\bf r}$$

2.  $E_{e-c} \neq \int n(\mathbf{r})V_{+}(\mathbf{r})d\mathbf{r}$ "usual" pseudopot. energy

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3.  $E_{\text{e-die}} = \int n(\mathbf{r}) V^{\text{die}}(\mathbf{r}) d\mathbf{r}$  could be bundled together with 1.

 $V_{+}(\mathbf{r})$ 

Coulombic potential due to the cores

$$V_{-}(\mathbf{r})$$

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 $V_{+}(\mathbf{r})$ 

Coulombic potential due to the cores

 $V_{-}(\mathbf{r})$ 

Coulombic potential due to the electrons  $V^{
m die}({f r})$ 

potential due induced charge (dielectric's response)

4. 
$$E_{\rm c-c} = \frac{1}{2} \sum_{A,B} \frac{Z_A Z_B}{R_{AB}}$$

Assuming open BC and  $\epsilon$ =1 between the cores

$$z(\mathbf{r}) = \sum_{A} Z_{A} \widetilde{\delta} \left( \mathbf{r} - \mathbf{R}_{A} \right)$$

charge density of spilled cores

 $n(\mathbf{r})$  charge density of electrons

 $q^{\rm die}({f r})$ 

induced charge density (not directly accessible from APBS)

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$$E_{\text{e-e}} = \frac{1}{2} \int n(\mathbf{r}) V_{-}(\mathbf{r}) d\mathbf{r}$$

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Coulombic potential due to the cores

$$V_{-}(\mathbf{r})$$

Coulombic potential due to the electrons  $V^{
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potential due induced charge (dielectric's response)

5. 
$$E_{\rm c-die} = \int z(\mathbf{r}) V^{\rm die}(\mathbf{r}) d\mathbf{r}$$

Assuming open BC and  $\epsilon$ =1 between the cores

4.  $E_{c-c}$ 

 $z(\mathbf{r}) = \sum_{A} Z_{A} \delta(\mathbf{r} - \mathbf{R}_{A})$ charge density of spilled cores

 $n(\mathbf{r})$ 

charge density of electrons



 $V_{+}(\mathbf{r})$ 

Coulombic potential due to the cores

$$V_{-}(\mathbf{r})$$

Coulombic potential due to the electrons  $V^{\text{die}}(\mathbf{r})$ 

 $q^{\text{die}}(\mathbf{r})$ 

induced charge density (not directly accessible from APBS)

1. 
$$E_{\text{e}-\text{e}} = \frac{1}{2} \int n(\mathbf{r}) V_{-}(\mathbf{r}) d\mathbf{r}$$

2.  $E_{e-c} \neq \int n(\mathbf{r})V_{+}(\mathbf{r})d\mathbf{r} V^{die}(\mathbf{r})$  can be only obtained numerically. ",usual" pseudopot. energy

3.  $E_{\text{e-die}} = \int n(\mathbf{r}) V^{\text{die}}(\mathbf{r}) d\mathbf{r}$ could be bundled together with 1.

#### We need to be able to calculate

 $\frac{\partial E_{\text{e-die}}}{\partial K^{\alpha\beta}}, \frac{\partial E_{\text{e-die}}}{\partial \omega_{\text{n}}(\mathbf{r})}, \frac{\partial E_{\alpha\beta}}{\partial k}$ 

$$\frac{\partial E_{\rm c-die}}{\partial K^{\alpha\beta}}, \frac{\partial E_{\rm c-die}}{\partial K^{\alpha\beta}}$$

$$rac{\delta E_{ ext{c-die}}}{\delta arphi_{ ext{\eta}}(\mathbf{r})}$$

potential due induced charge (dielectric's response)

5. 
$$E_{\rm c-die} = \int z(\mathbf{r}) V^{\rm die}(\mathbf{r}) d\mathbf{r}$$

Assuming open BC and  $\varepsilon$ =1 between the cores

4.  $E_{\rm c-c} = \frac{1}{2} \sum_{A,B} \frac{Z_A Z_B}{R_{AB}}$ 

#### Current state of affairs

- Ignore the gradient discrepancy until we think of something better.
- Backed up by the fact that V<sup>die</sup> is small compared to V<sub>-</sub> and V<sub>+</sub>.



Convergence -- standard ONETEP





#### Convergence -- standard ONETEP vs ONETEP+APBS





ONETEP+APBS, self-consistent (attempt 1)

2.16



Gaussian PCM	5.58	
ONETEP+APBS, frozen density, MP-corrected	7.60	
ONETEP+APBS, self-consistent (attempt 1)	67.96	

### **Further difficulties**

• The density-dependent cavity changes shape between iterations, because the density itself changes.

$$\frac{\delta E_{\rm es}}{\delta \rho(\mathbf{r})}(\mathbf{r}) = \phi(\mathbf{r}) + c(\nabla \phi(\mathbf{r}))^2 \frac{\delta \varepsilon(\mathbf{r})}{\delta \rho(\mathbf{r})}$$

#### Density-dependent cavity of Fattebert and Gygi



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 For now, to define the cavity, we are using the fixed density obtained from a converged calculation in vacuum.

#### **Further difficulties**

 APBS is memory-hungry. Treating large systems (proteins of several thousand atoms) at fine grid spacings requires 40-80 GB of memory.

<i>L</i> [Å]	RAM req'd	
10	80 MB	Assuming fine grid spacing of <i>d</i> =0.136 Å (KE cutoff ≈780 eV)
20	640 MB	
50	10 GB	
100	80 GB	
200	640 GB	

- APBS lacks a "proper" parallel implementation.
- There is a kludgy parallel mode for APBS, but it's Bad and Wrong.

### Conclusions

- Capturing solvation effects is important, particularly for biochemical applications.
- Implicit solvation models are computationally cheaper than explicitly modelling the solvent, but not necessarily easy to develop or to implement.
- Simple approach of correcting the *in vacuo* calculation with a postprocessing step of solving the PB equation for the frozen density appears to work surprisingly well.
- ... for the simplest molecules we've studied so far.
- Some clever maths will be required to further the advances towards a self-consistent approach.

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a Cheshire cat without the smile



just the smile, without the cat