Minimal-parameter implicit solvent model for large-scale DFT calculations

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with previous work by

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ONETEP Masterclass 2024

Solvation - why bother?

- Many important biochemical reactions occur in aqueous solution.
- Performing calculations in vacuo often leads to greatly inaccurate results. Especially sensitive properties and phenomena include, among others:
 - energy differences between molecular conformers [1],
 - rates of reactions [2],
 - tautomeric equilibria [1],
 - molecular (esp. protein-protein) associations [1,3],
 - protein structures [4],
 - ligand binding free energies [5].
- Thus, it is crucial to include the solvent environment in simulations of biological molecules.



Na⁺ and its 1st solvation shell



menthol molecule in water

Explicit solvent

- In *explicit solvent* methods we introduce the solvent in full atomic detail.
- (+) Accurate treatment of solute-solvent interactions.
- (-) Increase in system size, possibly by an order of magnitude.
- (-) Must average out instantaneous interactions (integrate out the degrees of freedom of solvent).
 - How to orient the solvent molecules?
 - How many configurations for averaging?
 - How to generate these configurations?



Phenol in explicit water Animation by Chris Pittock (priv. comm.)

Implicit solvent

- Only the solute is treated quantum-mechanically. We place it inside a suitably constructed dielectric cavity, whose inside is inaccessible to the solvent.
- The solvent is represented by an unstructured dielectric continuum. We only model its mean effect on the solute.
- (+) No solvent atoms (low cost).
- (+) Eliminates the costly sampling of solvent motions.
- (-) Simplified, mean-field model. All specific interactions between solute and solvent are lost.



General idea of *implicit solvent*. The solute is treated quantummechanically, while the solvent is represented by a dielectric continuum.







Elegant implicit solvent model

▶ We solve either the Generalized Poisson Equation:

 $\nabla\left(\varepsilon(\mathbf{r})\nabla\phi(\mathbf{r})\right) = -4\pi\rho(\mathbf{r})$

or the Poisson-Boltzmann equation (when there is electrolyte):

$$\nabla \cdot (\varepsilon(\mathbf{r}) \nabla \varphi(\mathbf{r})) = -4\pi \left[\rho(\mathbf{r}) + \sum_{i=1}^{p} z_i c_i(\mathbf{r}) \right] \qquad \begin{cases} z_i \}_{i=1}^{p} & \text{- charges of ion types} \\ \{c_i(\mathbf{r})\}_{i=1}^{p} & \text{- their concentration distributions} \end{cases}$$

... to get the electrostatic potential in solvent.

Two terms in free energy of solvation



- The electrostatic or polar term describes the response of the solvent to the charge distribution of the solute [7].
- It is the difference between the electrostatic energy

$$\frac{1}{2}\int\rho\left(\mathbf{r}\right)\phi\left(\mathbf{r}\right)d\mathbf{r}$$

in solvent and in vacuum.

- The nonpolar term accounts for
 - the entropic cost of forming a cavity within the solvent (cavitation energy),
 - for the van der Waals interaction of the solute with the solvent [4] (dispersion-repulsion energy).
- Difficult to describe rigorously.
- A widely used approach is to represent it as a linear function of the molecular surface area

[7]:
$$\Delta G_{npol} = \gamma A_{SA}$$
.

The procedure

- Perform a calculation in vacuo to obtain E_{vac} and the charge density in vacuum.
- Start a calculation in solvent, using the charge density in vacuum as initial guess.
- Generate the cavity basing on current charge density.
- Solve $\nabla \varepsilon(\mathbf{r}) \nabla \phi(\mathbf{r}) = -4\pi \rho(\mathbf{r})$ with a multigrid solver to obtain $\phi(\mathbf{r})$ in solvent. Use this in the electrostatic energy terms.
 - Repeat until convergence in solvent.



First achieve selfconsistency in vacuum.



... then in solution. NB how the density hardly changes at all in solvent.

In practice, it's as simple as that

! Turn on auto-solvation

is_auto_solvation T

! Define the permittivity of your solvent (default: water)

is_bulk_permittivity 78.54

! Define the Surface tension of your solvent (default: water)

is_solvent_surf_tension 0.07415 N/m

Calculation of free energy of	solvation:	hartree	kcal/mol
- Total energy in solvent:	(+)	-23.22902755568246	-14576.434548
- Total energy in vacuum:	(-)	-23.20990671966879	-14564.436043
- Total free energy of solvation	on:	-0.01912083601367	-11.998505

The above is a direct calculation of the free energy of solvation as a difference of the in-solvent and in-vacuum energies.

Components of polar term in f.e. of solvat	ion: hartree	kcal/mol
- Electrostatic:	-0.06759943752720	-42.419287
- Change in nonelectrostatic DFT terms:	0.04263084483500	26.751258
- Polar term in f.e. of solvation:	-0.02496859269221	-15.668028

The above is the calculation of the polar term to solvation, as a sum of the change in electrostatic energy between in-solvent and in-vacuum and the change in the remaining DFT terms.

Components of free energy of solvation:	hartree	kcal/mol
- Polar term in f.e. of solvation: (+)	-0.02496859269221	-15.668028
- Apolar (cavitation, dis., rep.): (+)	0.00584775667854	3.669523
- Total free energy of solvation:	-0.01912083601367	-11.998505

Finally, the total free energy of solvation is calculated as the sum of the polar and apolar terms calculated earlier. This is usually what you are after.



Read the manual!

ONETEP Documentation

7.1

Search docs

CONTENTS:

Starting with ONETEP

Working with GitHub

Ground State Calculation Setup

Using the Pseudoatomic Solver to Generate NGWFs

Conduction NGWF optimisation and optical absorption spectra

Finite-temperature DFT calculations using the Ensemble-DFT method

Running linear-scaling DFT calculations for metallic systems with the AQUA-FOE method

Density mixing (Kernel-DIIS)

Empirical Dispersion Correction

Using van der Waals Density Functionals

Realspace local pseudopotential in ONETEP

Solvent and Electrolyte Model

Overview of capabilities

E The models

A / Ground State Calculation Setup / Solvent and Electrolyte Model

View page source

Solvent and Electrolyte Model

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This manual pertains to ONETEP versions v6.0.0 and later. For older versions, see separate documentation on the ONETEP website. Major changes relative to v6.0.0:

- Soft-sphere model added in v6.1.1.8
- Surface Accessible Volume added in v6.1.3.0
- Conjugate gradient solver added in v6.1.3.6
- Self-consistent Continuum Solvation (SCCS) model added in v6.1.11.0
- Solvation forces in PBC added in v6.1.15.0
- Electrolyte forces added in v6.1.15.5
- Forces for soft-spheres solvation model added in v6.1.15.9

WARNING to users of v6.1.3.0 and later.

The method used to calculate the surface area of the dielectric cavity was changed in version 6.1.3.0. The surface area is used to calculate the $\Delta G_{\rm npol}$ component of the solvation. The new method is more mathematically consistent, but gives approximately 20% smaller values for the surface area. By default, we use the new method, which means the value of $\Delta G_{\rm solv}$ and may not



Do the tutorial!

Search docs

Introduction

Tutorial 1: Setting up Simple ONETEP Calculations

Tutorial 2: ASE ONETEP interface

Tutorial 3: Setting up Spin-Polarised Calculations

Tutorial 4: Geometry optimization

Tutorial 5: Analysis and visualization

Tutorial 6: Time-Dependent DFT

Tutorial 7: Spectral Function Unfolding

 Tutorial 8: Implicit solvation, visualisation and properties: Proteinligand free energy of binding for the T4 lysozyme

Introduction

 \boxplus Setting up the calculations

Evaluating the Outputs

Properties

Tutorial 9: DFT+U on strongly correlated magnetic materials: A case study on antiferromagnetic Hematite

Tutorial 10: Simulation cell relaxation

Tutorial 11: Electrified electrodeelectrolyte interfaces under potentiostatic control

Tutorial 12: Quantum embedding with (time-dependent) embedded mean-field theory: hydrogenation and excitations of pentacene

Tutorial 13: Electron Energy Loss Spectroscopy in ONETEP

Tutorial 8: Implicit solvation, visualisation and properties: Protein-ligand free energy of binding for the T4 lysozyme

Author: Lennart Gundelach, Jacek Dziedzic Date: June 2021 (revised June 2023)

Introduction

Protein-Ligand Free Energies of Binding

The binding free energy is a measure of the affinity of the process by which two molecules form a complex by non-covalent association. An example of this, of central importance in biology, is the binding of a ligand to a protein. Many methods to computationally approximate the binding free energies of protein-ligand interactions have been proposed with the ultimate goal of computationally predicting small molecule drug candidates which bind strongly to the protein of interest.

Quantum Mechanics in Binding Free Energies

A key limitation common to most computational methods of estimating binding free energies is the assumption of the validity of classical mechanics. The atoms and electrons that constitute biological molecules, like proteins, are, however, governed by the laws of quantum mechanics. Charge transfer, polarization and non-local interactions are not captured by traditional classical mechanical force-fields. Thus, a true description of protein-ligand binding requires a quantum mechanical (QM) treatment of the problem. In theory, a full, *ab-initio* QM approach would be system-independent, parameter-free and would describe the full spectrum of physical phenomena at work.

Unfortunately, high-level QM methods like coupled-cluster (CC) are prohibitively expensive and often have cubic or worse scaling with system size. Thus, even the ligands alone are often too large for routine calculations with these methods.



Results for small molecules

After we devised and implemented several corrections to the FGS model, we obtain very good accuracy with our model (MPSM).

Mean-square error:
 AMBER - 3.3 kcal/mol
 PCM - 4.9 kcal/mol
 FGS - 5.0 kcal/mol
 MPSM - 1.6 kcal/mol

AMBER - classical force field PCM - widely-used quantum approach FGS - model before our corrections MPSM - our model



Results for industrially-relevant molecules



		rms	\max	
Approach	XC functional	error	error	r
• this work ^{a}	PBE	3.8	8.3	0.83
this work ^{b}	PBE	4.1	9.1	0.83
$\circ \mathrm{PCM}$	PBE	10.9	23.3	0.53
◆SMD	M05-2X	3.4	14.5	0.87
AMBER	(classical)	5.1	19.9	0.77

Validation on 71 medium-size neutral molecules from

- Nicholls, Mobley, Guthrie, Chodera, Bayly, Cooper and Pande, "Predicting Small-Molecule Solvation Free Energies: An Informal Blind Test for Computational Chemistry", J. Med. Chem. 51 (2008).
- Guthrie, "A Blind Challenge for Computational Solvation Free Energies: Introduction and Overview", J. Phys. Chem. B 113 (2009).

Implicit solvation with thousands of atoms

- We applied our approach to the full (untruncated) human T4 lysozyme protein to study its free energies of binding to small ligands.
- Such system sizes (~2600 atoms) are out of reach of conventional DFT.

Molecule	$\Delta G_{ m lig,solv}^{ m SMD}$	$\Delta G^{\mathrm{exp}}_{\mathrm{lig,solv}}$	$\Delta G_{ m lig,solv}^{ m MM}$	$\Delta G_{\rm lig,solv}^{\rm QM}$	
Catechol	-9.3	-9.4 [54]	-20.9	-8.1	our protoin
3-Chlorophenol	-6.7	-	-9.9	-3.6	our protein,
2-Fluoroaniline	-4.5	_	-5.4	-3.2	with the
2-Aminophenol	-9.8	_	-13.9	-8.0	
2-Methylphenol	-6.3	-5.9 [55]	-9.0	-2.9	dielectric
1-Phenylsemicarbazide	-15.1	-	-16.2	-13.8	
Toluene	-1.3	-0.9 [55]	-1.4	1.4	cavity shown
Phenol	-6.7	-6.6 [55]	-9.7	-3.7	-
Relative hydration energies					
Catechol	-2.5	-2.8	-11.1	-4.3	
3-Chlorophenol	0.1	-	-0.1	0.1	
2-Fluoroaniline	2.3	-	4.4	0.6	
2-Aminophenol	-3.1	-	-4.1	-4.2	
2-Methylphenol	0.4	0.7	0.8	0.9	
1-Phenylsemicarbazide	-8.4	-	-6.4	-10.1	
Toluene	5.4	5.7	8.3	5.1	
Phenol (reference)	0.0	0.0	0.0	0.0	
Max error ^a			8.6	1.8	
RMS error ^a			3.4	1.2	

our protein,	
with the	
dielectric	

International Journal of International Journal of HEMISTRY



Implicit solvation with thousands of atoms

Hybrid functional (B3LYP) calculations on an aluminosilicate imogolite nanotube with 1416 atoms and implicit solvent, with near-complete basis set accuracy.





Other features not discussed here

- **Full support for PBCs and OBCs**. OBCs are used by default.
- Forces from all solvation terms:
 - geometry optimisation in solvent is possible,
 - ► MD in solvent is possible.
- Alternative solvation models available in ONETEP:
 - Fisicaro's soft-sphere model.
 - Andreussi's Self-Consistent Continuum-Solvation model.
- Solvent exclusion regions.
- Electrolyte.

Conclusions

- Building on the isodensity model of Fattebert and Gygi, we have developed a solvation model which:
 - ▶ is based on first principles,
 - has predictive power that is superior to classical models and to PCM,
 - ► has only two parameters (β , n_0), whose values have been optimized and which appear to be universal (do not depend on the solute),
 - b uses only two fundamental quantities (ϵ , γ), with clear physical interpretation, to describe the solvent.
- With our current implementation we can do implicit solvent calculations on systems with ~10⁴ atoms within a day.
- (!) ONETEP is free for academics. More info including tutorials and case studies: <u>www.onetep.org</u>.

Recommended reading

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